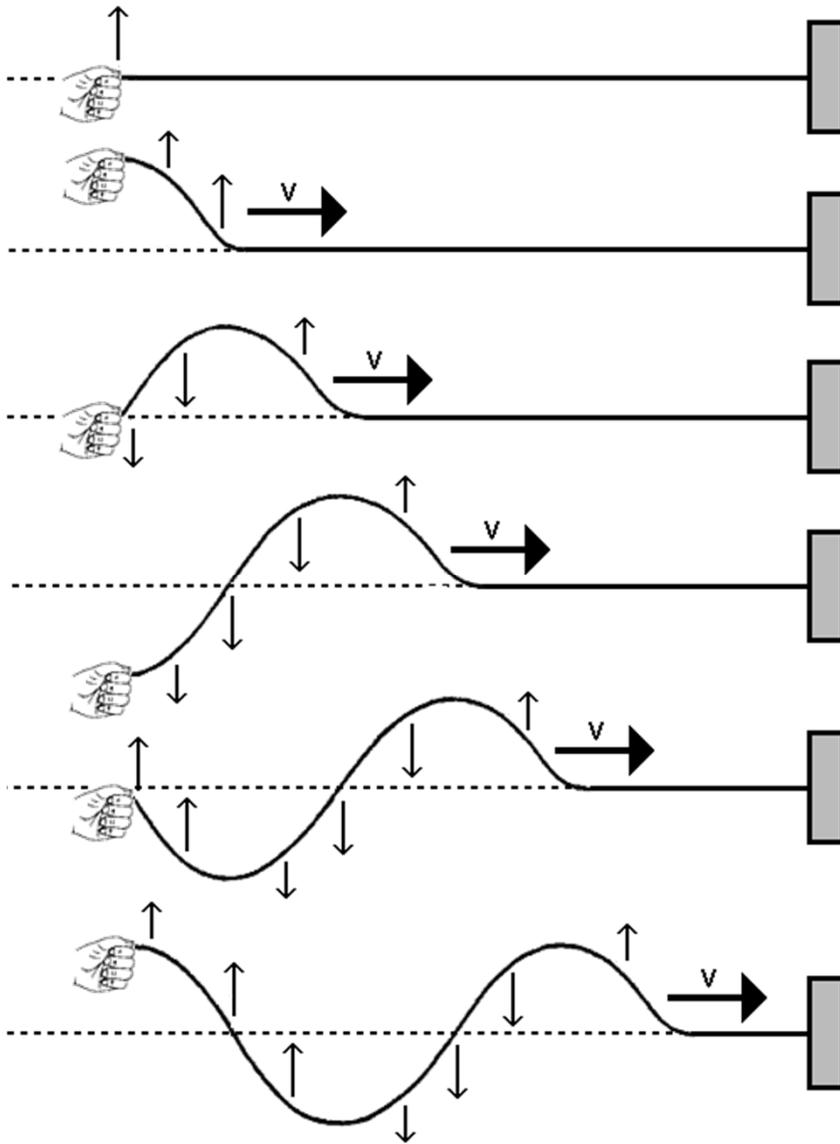
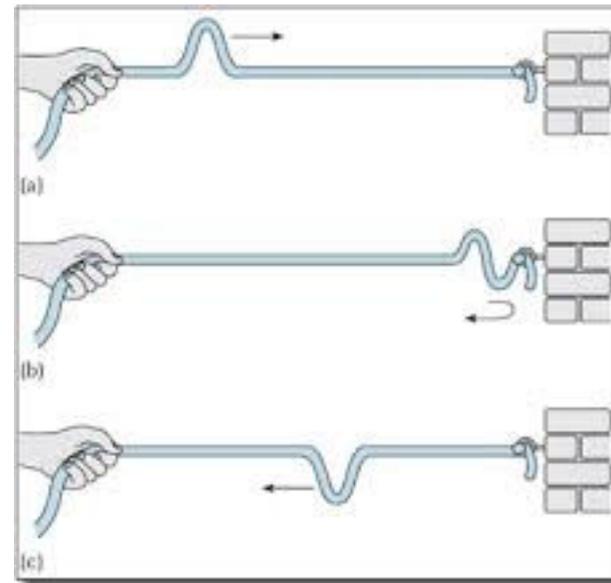


Traveling Waves

The one-dimensional (1D) case



A traveling wave is the propagation of motion (disturbance) in a medium.



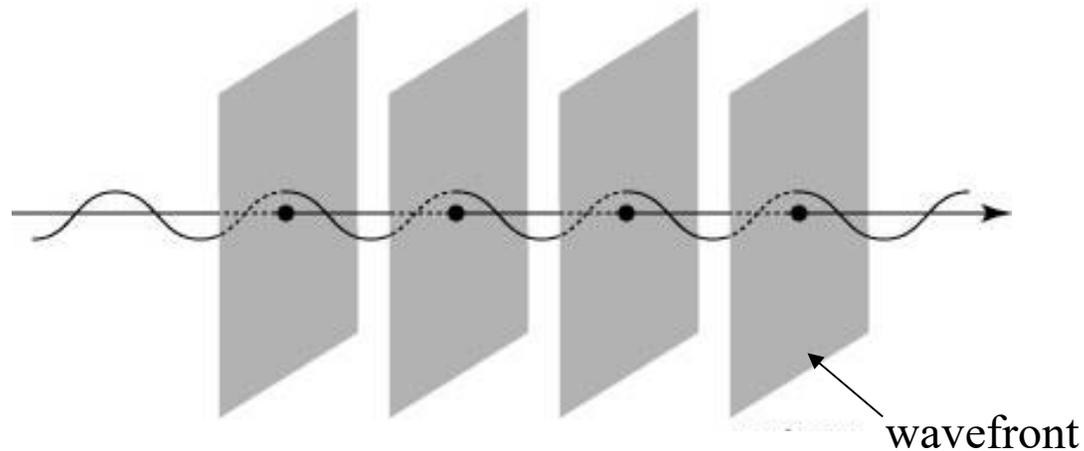
Reflection

Why is there reflection?

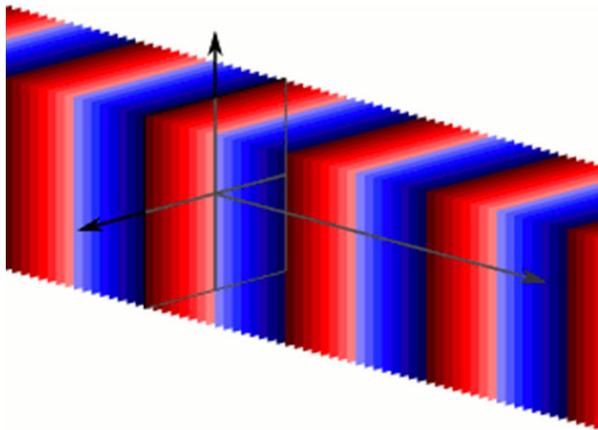
The "perturbation" propagates on.

Traveling Wave in Higher Dimensions

Plane waves in 3D



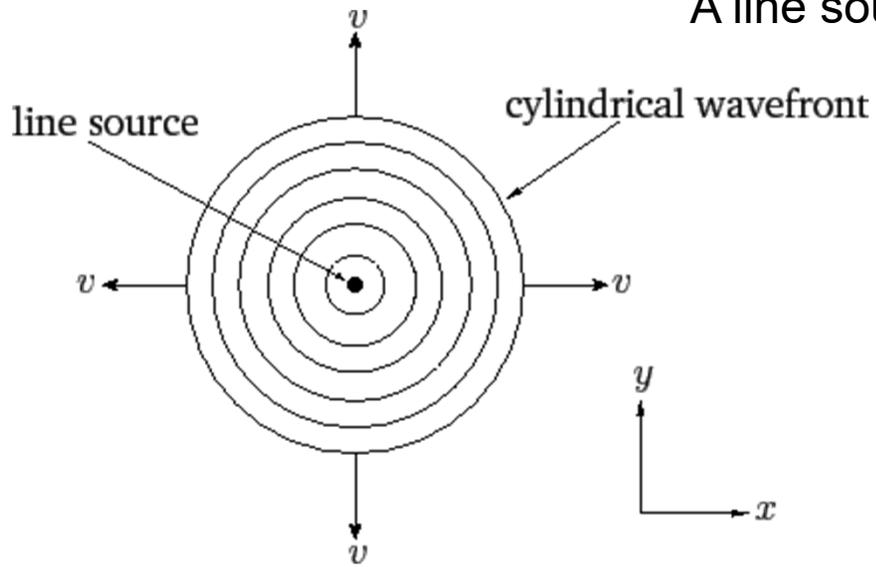
Example: sound waves



The plane wave is a 1D wave in 3D space:
No variation in the 2D plane of a wave front.
A wave front is a surface of equal phase.

Watch animation: http://en.wikipedia.org/wiki/Plane_wave

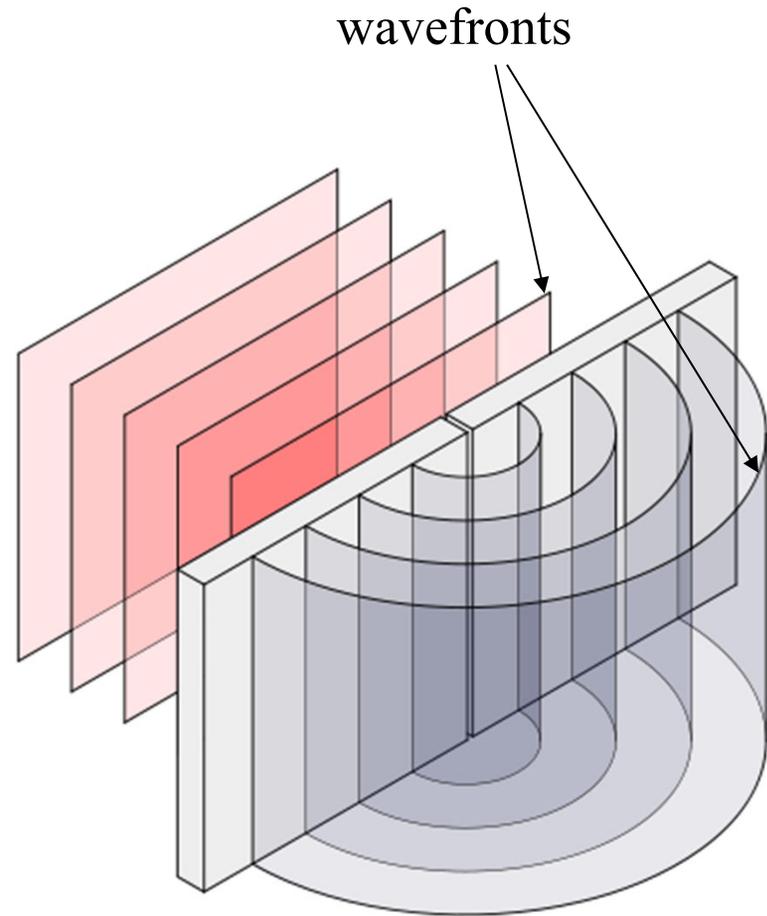
A line source makes a cylindrical wave.



Cylindrical wave (3D; top view)



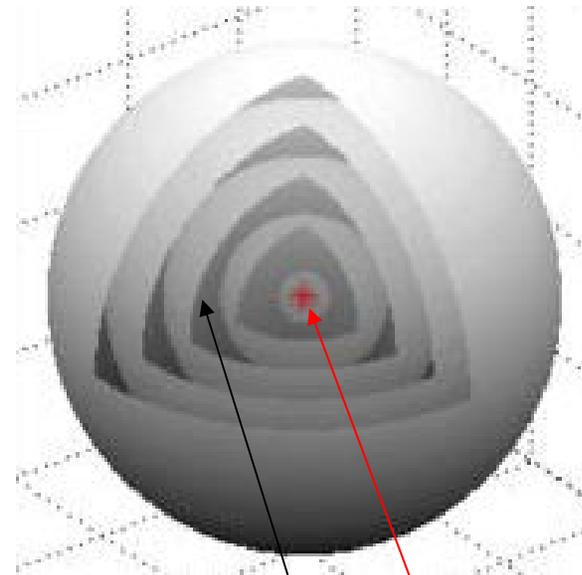
Water surface wave (2D)
(Circular wave)



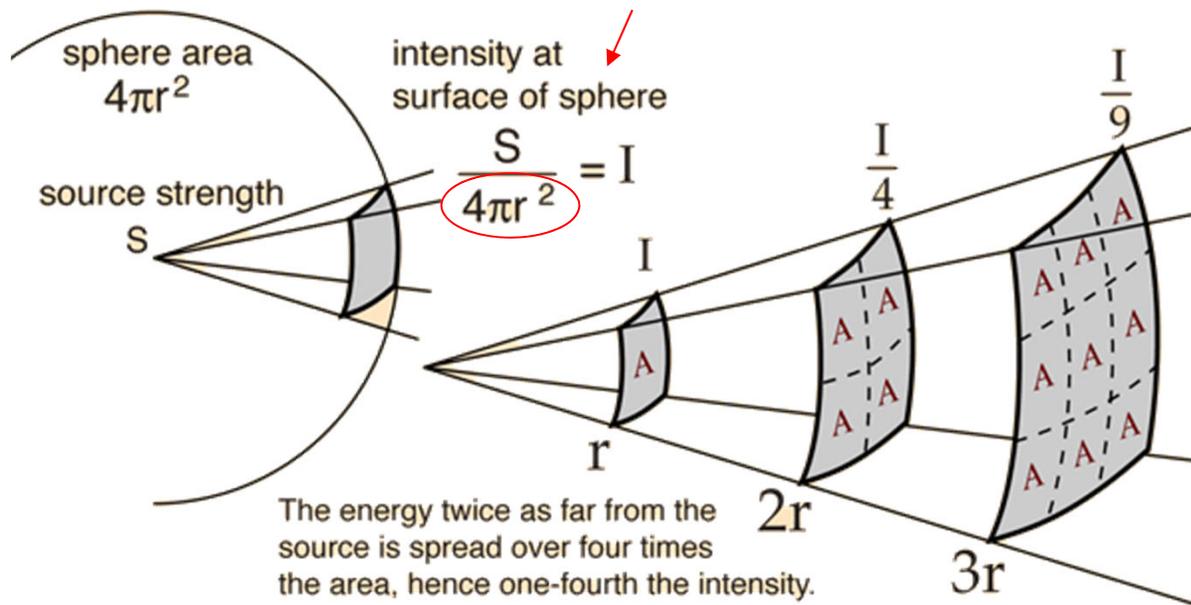
Make a cylindrical wave from a
plane wave

A point source makes a spherical wave, the wave front of which are spherical surfaces.

Intensity is energy carried per time per area, i.e., power delivered per area



Conservation of energy



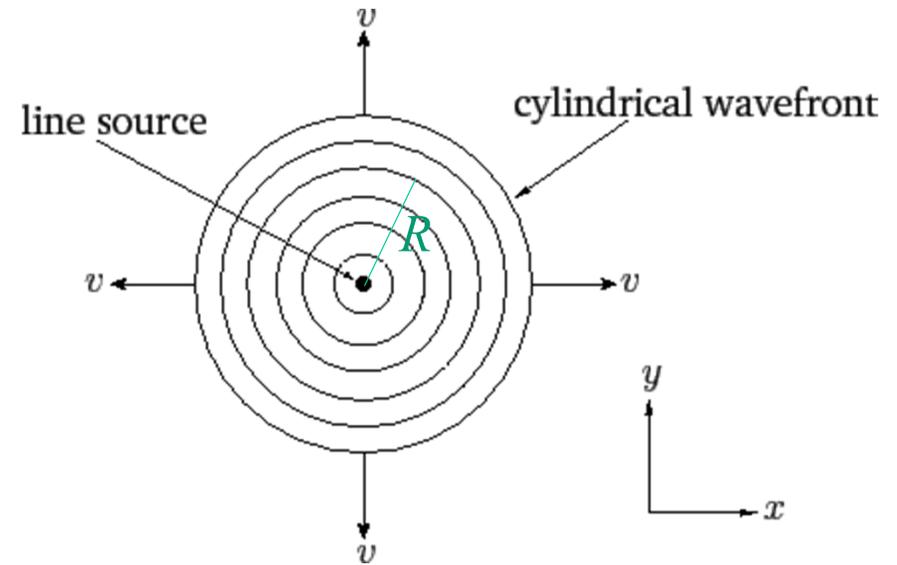
Point source

wavefronts

Intensity $\propto 1/r^2$
 Amplitude $\propto 1/r$

How does the intensity & amplitude of a cylindrical wave depend on R ?

$$R^2 = x^2 + y^2$$



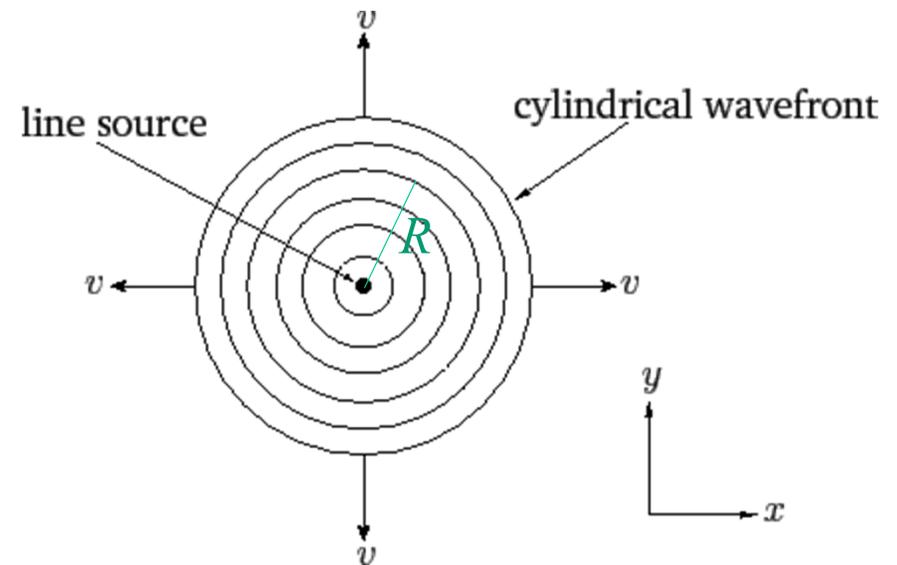
How does the intensity & amplitude of a cylindrical wave depend on R ?

$$R^2 = x^2 + y^2$$

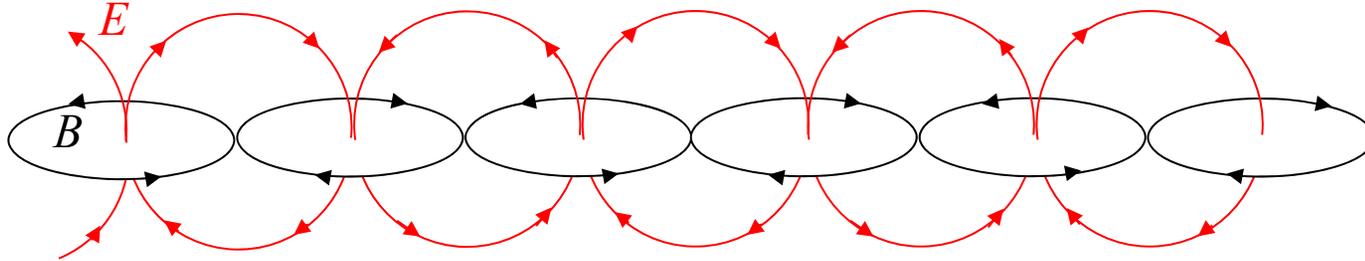
Circumference $\propto R$

Therefore,

$\begin{aligned} \text{Intensity} &\propto 1/R \\ \text{Amplitude} &\propto 1/(R^{1/2}) \end{aligned}$
--



Electromagnetic Wave



Somehow start with a changing electric field E , say $E \propto \sin \omega t$

The changing electric field induces a magnetic field, $B \propto \frac{\partial E}{\partial t} \propto \cos \omega t$

If the induced magnetic field is changing with time, it will in turn induce an electric field

$$E \propto -\frac{\partial B}{\partial t} \propto \sin \omega t$$

Notice that $\frac{d}{dt} \cos \omega t = -\omega \sin \omega t$

Negative signs cancel

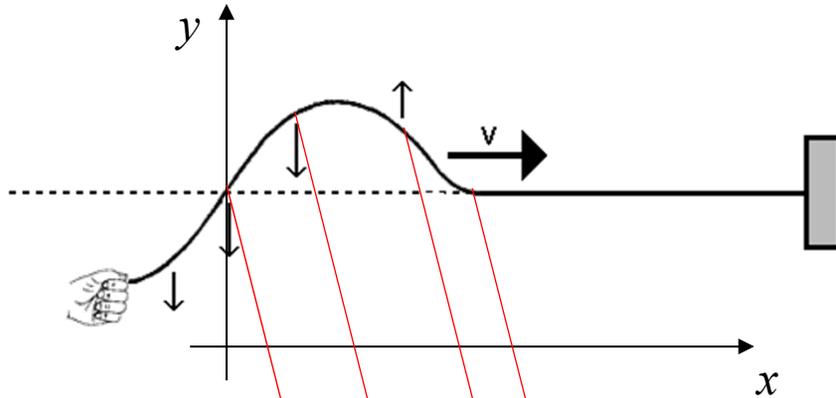
And so on and so on....

Just as the mechanical wave on a string.

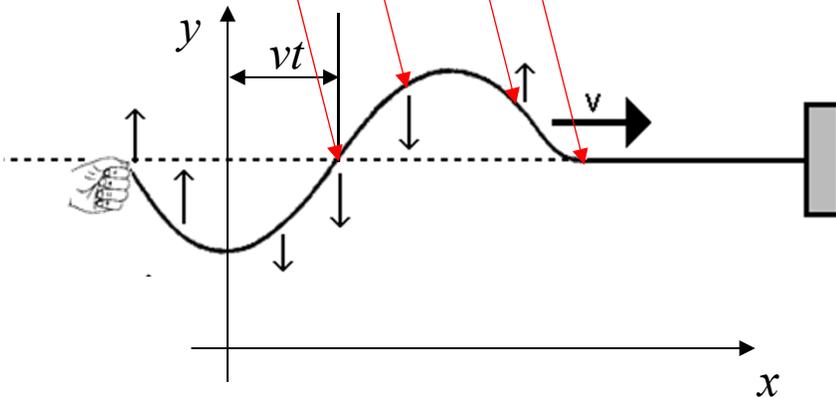
Note: The picture is NOT of a plane wave, but depicts a wave emitted by a source of limited size. Will revisit this picture later.

Mathematical Expression of the Traveling Wave

A traveling wave is the propagation of motion (disturbance) in a medium.



At time 0,
 $y = f(x)$



At time t ,
 $y = f(x - vt)$

The wave form shifts vt in time t .

This is the general expression of traveling waves.

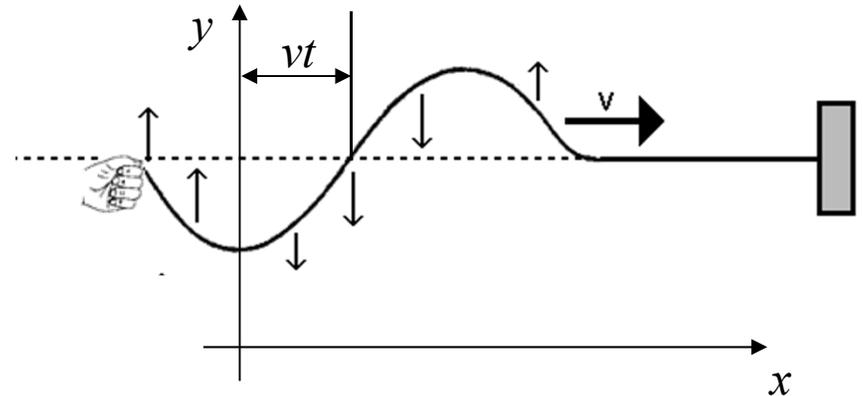
Questions:

What kind of wave does $y = f(x + vt)$ stand for?

What about $y = f(vt - x)$?

What about $y = f(vt - x)$?

$$f(vt - x) = f[-(x - vt)]$$



Define $f(-x)$ your “new f ”, or $g(x) \equiv f(-x)$, so it’s the same wave!

So, which way should I go? $f(x - vt)$ or $g(vt - x)$?

To your convenience!

No big deal. But this affects how we define our “sign conventions.”

People in different disciplines use different conventions.

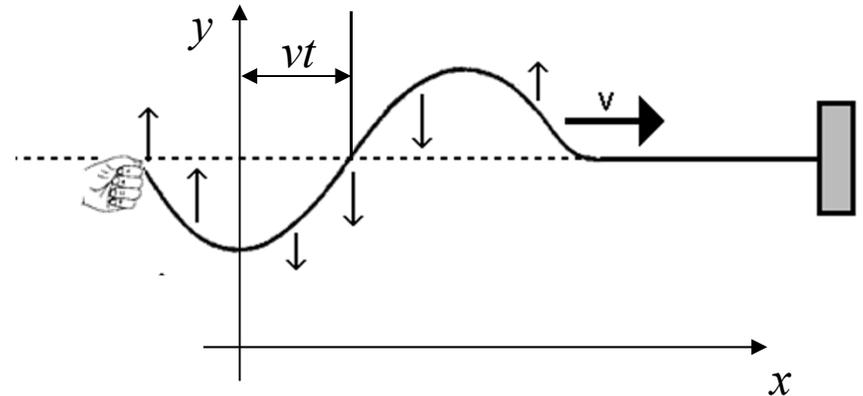
If you are more concerned about seeing a waveform on an oscilloscope, you like $g(vt - x)$ better.

If you are more concerned about the spatial distribution of things, you like $f(x - vt)$ better.

We will talk later about how this choice affects the ways to write the “same” (but apparently different) equations in different disciplines (EE vs. physics).

What about $y = f(vt - x)$?

$$\begin{aligned} f(vt - x) \\ = f[v(t - x/v)] \end{aligned}$$



At any x , you have a time-delayed version of $f(vt)$.

The time delay is simply the time for the wave to travel a distance x , i.e., x/v .

For a single-wavelength, sinusoidal wave, this is always true.

Because the wave travels at just one speed, v .

But, a general wave has components of different wavelengths/frequencies.

The speeds of the different components may be different.



I am already talking about an important concept.

Then, a distance x later, the “waveform” in time (as you see with an oscilloscope) will change.

This is called “dispersion”

Let's now look at the special case of the sinusoidal wave

$$\begin{aligned}y(x, t) &= A \cos(\omega t - \beta x + \phi_0) \\&= A \cos\left[\beta\left(\frac{\omega}{\beta}t - x + \frac{\phi_0}{\beta}\right)\right] \\&= A \cos[\beta(vt - x) + \phi_0] \\&= f(x - vt) = g(vt - x)\end{aligned}$$

ω : angular frequency

β : propagation constant

Phase as a function of x

You can write this function, or group the terms, in so many ways.
It's just about how you view them

Watch Wikipedia animation:

https://en.wikipedia.org/wiki/Wave#/media/File:Simple_harmonic_motion_animation.gif

Next, we try to understand β

$$y(x, t) = A \cos(\omega t - \beta x + \phi_0)$$

$$= A \cos\left[\beta\left(\frac{\omega}{\beta}t - x + \frac{\phi_0}{\beta}\right)\right]$$

$$= A \cos[\beta(vt - x) + \phi_0]$$

$$= f(x - vt) = g(vt - x)$$

$$v = \frac{\omega}{\beta} = \frac{2\pi f}{2\pi/\lambda} = f\lambda$$

$$= \frac{2\pi/T}{2\pi/\lambda} = \frac{\lambda}{T}$$

One wavelength traveled in one period.

λ is the "spatial period"; $\frac{1}{\lambda}$ is the "spatial frequency".

And, β is the spatial equivalent of ω .

Call it the wave vector or propagation constant.

For the free space (i.e. vacuum), $v = c = \frac{\omega}{\beta}$, or $\omega = c\beta$.

No dispersion.
(c is a constant)

$$y(x, t) = A \cos(\omega t - \beta x + \phi_0)$$

$$= A \cos\left[\beta\left(\frac{\omega}{\beta}t - x + \frac{\phi_0}{\beta}\right)\right]$$

$$= A \cos[\beta(vt - x) + \phi_0]$$

$$= f(x - vt) = g(vt - x)$$

$$v = \frac{\omega}{\beta} = \frac{2\pi f}{2\pi/\lambda} = f\lambda$$

$$= \frac{2\pi/T}{2\pi/\lambda} = \frac{\lambda}{T}$$

One wavelength traveled in one period.

For the free space (i.e. vacuum), $v = c = \frac{\omega}{\beta}$, or $\omega = c\beta$.

No dispersion (c is a constant)

The relation between β and ω for a wave traveling in a medium is a material property of the medium. We call it the “dispersion relation” or just “dispersion.”

Recall that we used the term to describe a phenomenon. Related.

In general the dispersion relation is not perfectly linear.

$\frac{\omega}{\beta} \equiv v$ is not a constant. We call it the phase velocity, v_p .

Thus the dispersion!

$$\begin{aligned}
 y(x, t) &= A \cos(\omega t - \beta x + \phi_0) \\
 &= A \cos\left[\beta\left(\frac{\omega}{\beta}t - x + \frac{\phi_0}{\beta}\right)\right] \\
 &= A \cos\left[\beta(vt - x) + \phi_0\right]
 \end{aligned}$$

ϕ_0 is the reference phase (the wave's phase with time and space set to zero)

$$\begin{aligned}
 y(x, t) &= A \cos\left[\omega\left(t + \frac{\phi_0}{\omega}\right) - \beta x\right] \\
 &= A \cos\left[\omega t - \beta\left(x - \frac{\phi_0}{\beta}\right)\right]
 \end{aligned}$$

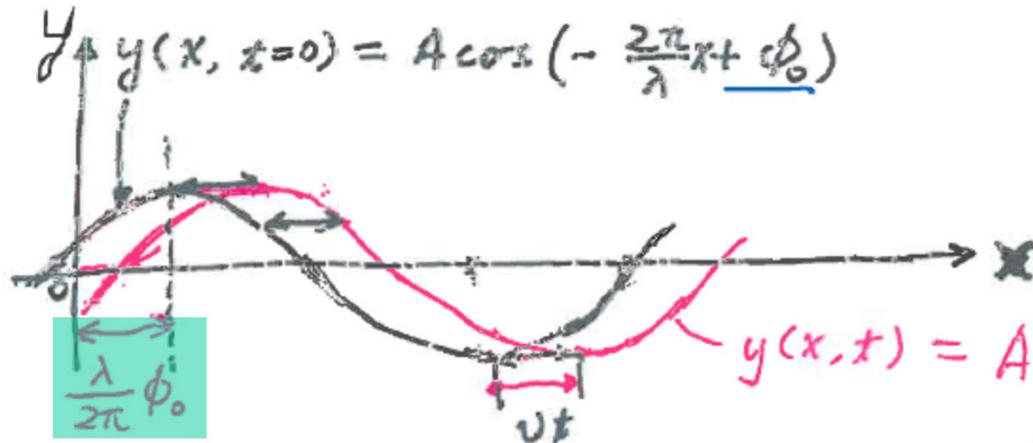
ϕ_0 viewed as shift in time

ϕ_0 viewed as shift in position

Two ways to look at this.
Two ways to group the terms.

$$y(x, t) = A \cos(\omega t - \beta x + \phi_0) = A \cos\left(\frac{2\pi}{T} t - \frac{2\pi}{\lambda} x + \phi_0\right)$$

$$y(x, t=0) = A \cos\left(-\frac{2\pi}{\lambda} x + \phi_0\right)$$



"spatial period"

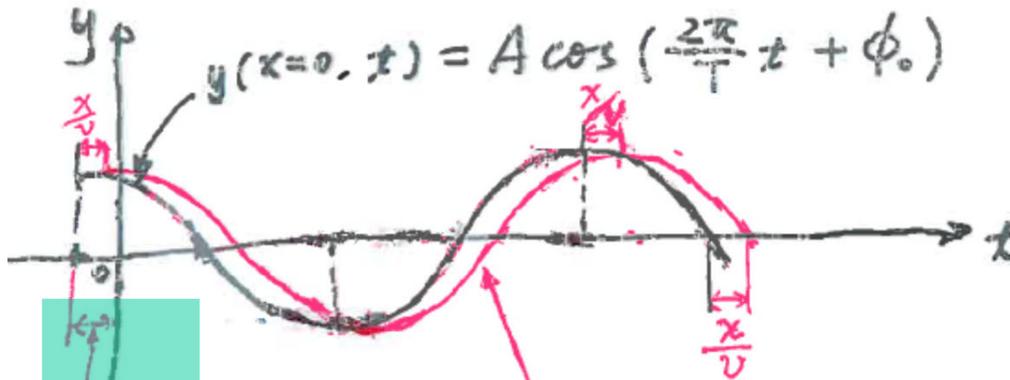
Take snapshots at different times

$$y(x, t) = A \cos\left[\left(-\frac{2\pi}{\lambda} x + \phi_0\right) + \frac{2\pi}{T} t\right]$$

$$= A \cos\left[-\frac{2\pi}{\lambda} x + \left(\phi_0 + \frac{2\pi}{T} t\right)\right]$$

Convert phase to distance

For the same wave,



Measure waveforms at different locations

$$y(x, t) = A \cos\left(\frac{2\pi}{T} t + \phi_0 - \frac{2\pi}{\lambda} x\right) = A \cos\left[\frac{2\pi}{T} t + \left(\phi_0 - \frac{2\pi}{\lambda} x\right)\right]$$

$$\frac{T}{2\pi} \phi_0$$

Convert phase to time

Now, you can work on HW1, P1 – P7.

Both the Homework & Answer sheet are online.

Review class notes and read the textbook, **then** do the homework, **then** check answers.

Waves carry information.

How much information does a sinusoidal wave carry?

Why do we study sinusoidal waves?

We want the wave to carry the “undistorted” information after it travels a distance x to reach us.

We want its snapshots in space to be the same as at $t = 0$.

We want its waveform in time to be the same as at $x = 0$.

Recall that we talked about dispersion.

In most cases, the dispersion is not too bad.

The $\omega(\beta)$ is only slightly nonlinear.

The “signal” or “wave packet” or “envelope” travels at a different speed than v_p , which is different for different frequencies anyway.

That speed is the “group velocity” v_g .

Run the extra mile:

Find out the expression for v_g , given the dispersion $\omega(\beta)$. Derive it.

You’ll have a deep understanding about wave propagation.

Attenuation

In some cases, the amplitude A decreases as the wave propagates.

For many types of waves, the power density $\propto A^2$

For unit distance traveled, a certain fraction of power density is lost.

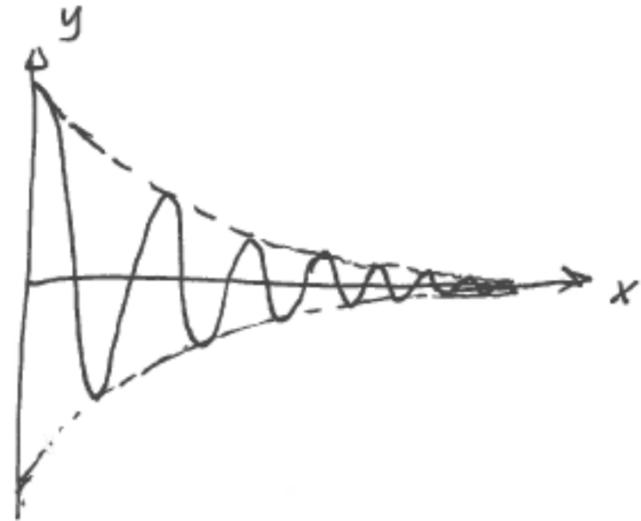
$$\frac{d A^2(x)}{A^2(x)} = -2\alpha dx$$

$$\frac{d A^2(x)}{dx} = -2\alpha A^2(x)$$

$$A^2(x) = A^2(0) e^{-2\alpha x}$$

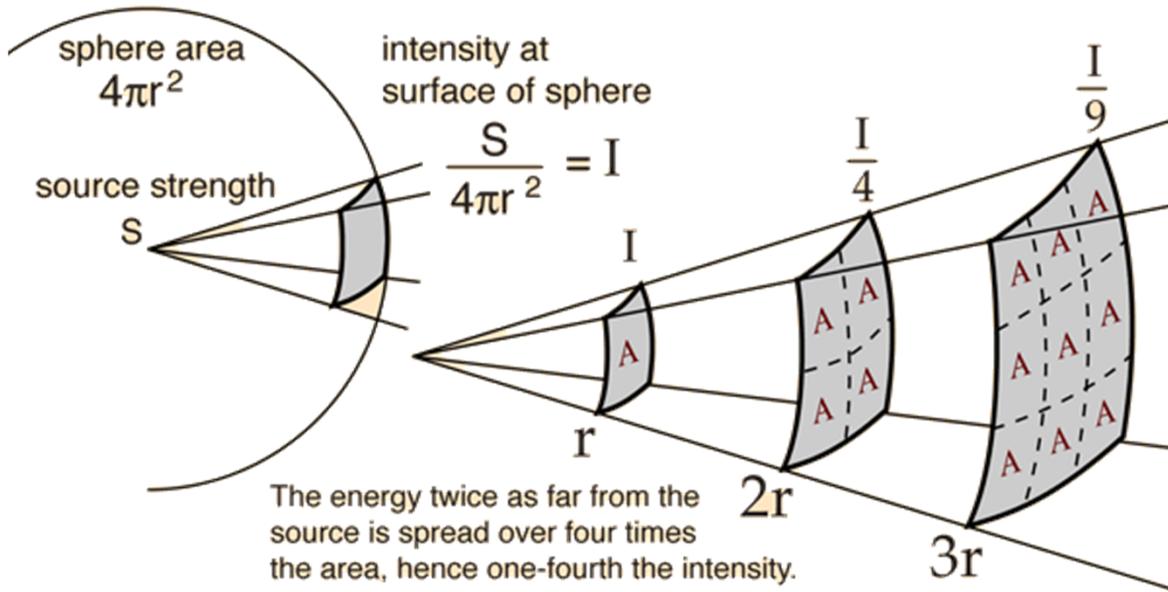
$$A(x) = A(0) e^{-\alpha x}$$

$$y(x, t) = A_0 e^{-\alpha x} \cos(\omega t - \beta x + \phi_0)$$



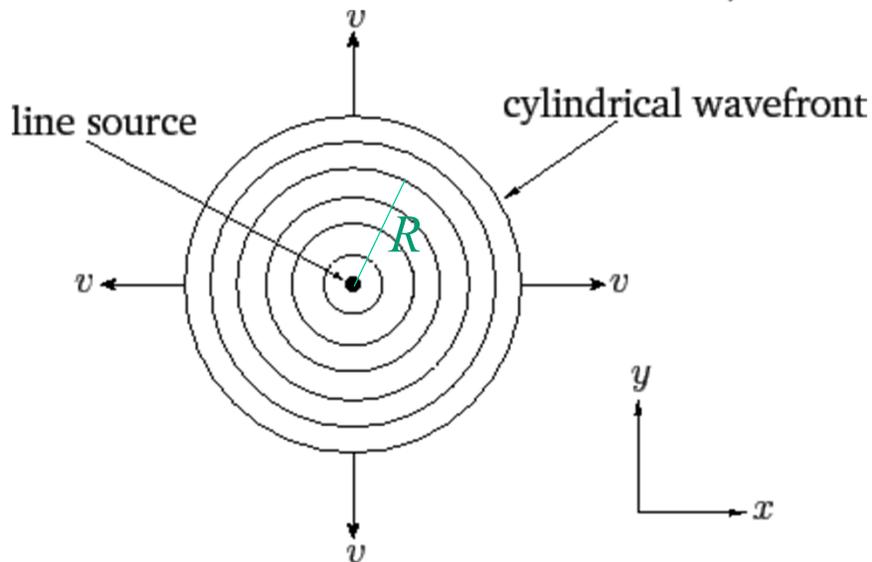
$$y(x, t) = A_0 e^{-\alpha x} \cos(\omega t - \beta x + \phi_0)$$

The exponential attenuation of the plane wave is due to loss, not to be confused with the amplitude reduction due to “spreading”.



$$y(r, t) = A_0(1/r)\cos(\omega t + \beta r + \phi_0)$$

Intensity $\propto 1/r^2$
Amplitude $\propto 1/r$



$$y(r, t) = A_0(1/R^{1/2})\cos(\omega t + \beta R + \phi_0)$$

Intensity $\propto 1/R$
Amplitude $\propto 1/(R^{1/2})$

Just to mention another mechanism that weakens a propagating wave: scattering

Finish HW1:P1 through P9.

Both the Homework & Answer sheet are online.

Review class notes and read the textbook, **then** do the homework, **then** check answers.