# **How to Study and Distinguish Forced Oscillations**

#### **Lei CHEN**

PhD, Associate Professor Department of Electrical Engineering Tsinghua University chenlei08@tsinghua.edu.cn





# **Natural and Forced Oscillations**

- Two types of oscillations are widely observed
- Natural/Free oscillation Oscillations due to undamped system modes
- Forced oscillation Oscillations from periodic sources external to the system

Natural and Forced Oscillations		
Two types of oscillations are widely observed		
Natural/Free oscillation - Oscillations due to		
Forced oscillation - Oscillations from periodic sources		
external to the system		
$\dot{x} = Ax$	$x(t) = \sum_{i=1}^{n} \varphi_i c_i e^{\lambda_i t}$	Natural
$\dot{x} = Ax + f(t)$	$f(t)$ is periodic	Forced

\n**IEEE**





# **Oscillation Type Distinguishing**

- Why should we distinguish oscillation type?
	- Different control measures for different oscillation types
	- Natural oscillation: Increase the damping ratio of the critical mode and the oscillation will decay
	- Forced oscillation: Remove the external disturbance
- Why is this problem difficult?
	- The approach should be measurement-based, or it is not online applicable
	- Both oscillation types show sustained oscillation with constant amplitude in steady state
	- Actual oscillation waveforms are complicated

















# **How to Study Forced Oscillation**

- Natural oscillation
	- Determined by system features
	- Eigenvalue analysis or modal analysis
	- Time-domain simulation
- Forced oscillation
	- Determined by both external disturbances and system features
	- Time-domain simulation method is applicable, but not capable of analytic and quantitative analysis
	- Extended modal analysis





## **Extended Modal Analysis**

Linearized system  $\dot{x} = Ax + Bu$ 

 $r=1$ 

*M*odal transformation  $x = \mathbf{\Phi} z$ al transformation  $x =$ <br>
bupled system  $\dot{z} =$ <br>
ural oscillation  $\begin{cases} z_r \\ z_r^* \end{cases}$ <br>  $u = 0$ 

Uncoupled system

$$
\dot{z}=Az+\boldsymbol{\Phi}^{-1}\boldsymbol{B}u
$$

Natural oscillation

$$
\begin{cases} z_r = z_{r0} e^{\lambda_r t} = \mathbf{\mathbf{\mathcal{Y}}}_r^T \mathbf{x}_0 e^{\lambda_r t} \\ z_r^* = z_{r0}^* e^{\lambda_r^* t} = \mathbf{\mathcal{Y}}_r^{*T} \mathbf{x}_0 e^{\lambda_r^* t} \end{cases}
$$



$$
x_i(t) = \sum_{r=1}^{n-1} |\phi_{ir}||z_{r0}| e^{-\alpha_r t} \left[ e^{j(\omega_{dr}t + \gamma_{ir} + \theta_r)} + e^{-j(\omega_{dr}t + \gamma_{ir} + \theta_r)} \right]
$$
  
=  $2 \sum_{r=1}^{n-1} |\phi_{ir}||z_{r0}| e^{-\alpha_r t} \cos(\omega_{dr}t + \gamma_{ir} + \theta_r)$ 

$$
\left|\phi_{ir}\right| \angle \gamma_{ir}
$$

Mode shape







#### **Extended Modal Analysis**

Forced oscillation

$$
\begin{bmatrix} \dot{z} \\ \dot{z}^* \end{bmatrix} = \begin{bmatrix} A \\ A^* \end{bmatrix} \begin{bmatrix} z \\ z^* \end{bmatrix} + \boldsymbol{\Phi}^{-1} \boldsymbol{B} \boldsymbol{u}
$$

Assumption: only one sinusoidal disturbance  $\bm{B} \bm{u} = \Delta \bm{P}_T = \Delta \bm{P}_{Tm} \sin \omega t$ 

Focus on steady state response

$$
\Delta \bm{P}_T = \Delta \tilde{\bm{P}}_{Tm} e^{\text{j} \omega t}
$$

$$
\begin{cases}\nz_r = \frac{\tilde{\mathbf{F}}_r^T \Delta \tilde{\mathbf{P}}_{Tm}}{j \omega - \tilde{\lambda}_r} e^{j \omega t} \\
z_r^* = \frac{\tilde{\mathbf{F}}_r^{*T} \Delta \tilde{\mathbf{P}}_{Tm}}{j \omega - \tilde{\lambda}_r^*} e^{j \omega t}\n\end{cases}
$$

$$
\Delta P_T = \Delta \tilde{P}_{Tm} e^{j\omega t}
$$
\nComplex phasor representation\n
$$
\begin{bmatrix}\nz^*_{r} = \frac{\tilde{\mathbf{Y}}_r^{*T} \Delta \tilde{P}_{Tm}}{j\omega - \tilde{\lambda}_r^{*}} e^{j\omega t} \\
\bar{\lambda}_r^{*} = \frac{\tilde{\mathbf{Y}}_r^{*T} \Delta \tilde{P}_{Tm}}{j\omega - \tilde{\lambda}_r^{*}} e^{j\omega t}\n\end{bmatrix}
$$
\n
$$
x(t) = \sum_{r=1}^{n} \underbrace{\tilde{\mathbf{\Phi}}_r^{*} z_r}_{=1} + \sum_{r=1}^{n} \underbrace{\tilde{\mathbf{\Phi}}_r^{*} z_r^{*}}_{(j\omega - \tilde{\lambda}_r^{*}) - (\tilde{\mathbf{\Phi}}_r^{*} \tilde{\mathbf{\Psi}}_r^{*} \tilde{\lambda}_r^{*} + \tilde{\mathbf{\Phi}}_r^{*} \tilde{\mathbf{\Psi}}_r^{*T} \tilde{\lambda}_r)}_{(j\omega - \tilde{\lambda}_r^{*})} \Delta \tilde{P}_{Tm} e^{j\omega t}
$$





# **Steady State Response of Forced Oscillation**

The rth mode in the *i*th state variable 
$$
x_{i,r}(t) = B_{i,r} \sin(\omega t - \varphi_{i,r})
$$
  
\n
$$
B_{i,r} = \sqrt{\frac{(a/\omega_m^2)^2 + (v_r b/\omega_m)^2}{(1 - v_r^2)^2 + (2\zeta_r v_r)^2}} |\tilde{p}_t| \qquad \varphi_{i,r} = \arctan \frac{2\zeta_r v_r a - (1 - v_r^2)v_r b\omega_m}{a(1 - v_r^2) + 2\zeta_r v_r^2 b\omega_m}
$$
\n
$$
a = -(\tilde{\lambda}_r^* \tilde{\phi}_r \tilde{\psi}_r + \tilde{\lambda}_r \tilde{\phi}_r^* \tilde{\psi}_r^*) \qquad \zeta_r \qquad \omega_m = |\lambda_r| \qquad v_r = \omega/\omega_m
$$
\n
$$
b = \tilde{\phi}_r \tilde{\psi}_r + \tilde{\phi}_r^* \tilde{\psi}_r^* \qquad \text{Damping ratio} \qquad \text{Natural frequency} \qquad \text{Frequency ratio}
$$
\nFrequency ratio inducing the largest amplitude  $v_r^2 = \sqrt{1 + 2(1 - 2\zeta_r^2)(\frac{a}{b\omega_{nr}})^2 + (\frac{a}{b\omega_{nr}})^4} - (\frac{a}{b\omega_{nr}})^2$   
\n $\zeta_r \approx 0 \qquad v_r \approx 1$   
\nPoory damped Resonance

Power

Е

Resonance - When the frequency of the external disturbance is close to the frequency of a poorly damped mode, the system oscillates with a large amplitude, which is often larger than that of the disturbance



#### Easy to Understand





## **Steady State Response at Resonance**

When resonance with a poorly damped mode, the mode dominates the response

$$
x_i(t) \approx \frac{\left|\tilde{\phi}_{ir}\right| |\tilde{\psi}_{lr}| |\tilde{p}_l|}{\zeta_r \omega_{nr}} \sin\left(\omega_{nr} t + \gamma_{ir} + \sigma_{lr}\right)
$$

$$
\zeta_r \qquad \tilde{\phi}_{ir} = \left| \tilde{\phi}_{ir} \right| e^{j\gamma_{ir}} \qquad \tilde{\psi}_{lr} = \left| \tilde{\psi}_{lr} \right| e^{j\sigma_{lr}}
$$
\nDamping ratio

\nRight eigenvector

\nLeft eigenvector

\ni

\nJ

\nState variable

\nLocation of disturbance





# **Oscillation Amplitude**

**Solution Arithmetic**  
\n
$$
x_{i}(t) \approx \frac{\left|\tilde{\phi}_{ir}\right| \left|\tilde{\nu}_{lr}\right| \left|\tilde{p}_{l}\right|}{\zeta_{r} \omega_{nr}} \sin\left(\omega_{nr} t + \gamma_{ir} + \sigma_{lr}\right)
$$
\nAmplitude of disturbance

\n
$$
\left|\tilde{P}\right|
$$

Oscillation amplitude is affected by

Damping ratio of mode Location of disturbance  $\left|\tilde{p}_{\overline{l}}\right|$  $\zeta_r$ 

- Larger amplitude of disturbance induces larger amplitude of oscillation
- Larger  $|\tilde{\psi}_{lr}|$ , which means the location of disturbance has stronger controllability on the mode, induces larger amplitude of oscillation
- Smaller damping ratio induces larger amplitude of oscillation





# **Mode Shape**

$$
x_i(t) \approx \frac{\left|\tilde{\phi}_{ir}\right| |\tilde{\psi}_{lr}| |\tilde{p}_l|}{\zeta_r \omega_{nr}} \sin\left(\omega_{nr} t + \gamma_{ir} + \sigma_{lr}\right)
$$

- Relative amplitude and phase of different state variables is also determined by the right eigenvector  $\left|\widetilde{\phi}_{ir}\right| \angle {\gamma}_{ir}$ **Mode Shape**<br>  $+\gamma_{ir} + \sigma_{lr}$ <br>
whase of different state variables is also detern<br>  $\left|\tilde{\phi}_{ir}\right| \angle \gamma_{ir}$ <br>
wh
- Same as natural oscillation
- Not affected by the location of disturbance

At resonance, the mode shape of forced oscillation converges to system mode shape, which brings difficulties in oscillation type distinguishing and source location





## **Simulations**



Inter-area mode Right eigenvector



 $\zeta = 2.15\%$ 





# **Simulations**









Е

# **How to Distinguish Forced Oscillation**

- Distinguishing natural and forced oscillations based on system measurements is critical for control measure decisions
- In steady state, both natural and forced oscillations show sustained oscillations with nearly constant amplitude
	- A damped oscillation is obviously natural oscillation Not Critical
	- Oscillation with negative damping will converge to constant-amplitude oscillation due to nonlinearities such as saturations and limits in actual systems
- Mode shapes are similar when resonance
- Not easy to distinguish





# **Fundamental Differences**

- Influence factor
	- Natural oscillation system features
	- Forced oscillation system features + external disturbances
	- The steady state waveform of natural oscillation is mainly sinusoidal
	- If the external disturbance is non-sinusoid, the forced oscillation waveform will also deviate a lot from sinusoid
	- An obvious non-sinusoidal waveform is a sufficient but unnecessary indicator of forced oscillation
- Intrinsic system damping
	- Natural oscillation zero or negative
	- Forced oscillation positive
	- How to obtain the intrinsic damping from its outward performance?





# **1. Harmonic Content of Steady State Waveform**

- Obvious non-sinusoidal waveform in steady state is a sufficient but unnecessary condition of forced oscillation 2 3 **Harmonic Content of**<br>bvious non-sinusoidal was<br>afficient but unnecessary cor<br>armonic content<br>- Harmonic content higher than a<br>sinusoidal waveform, and forced<br> $m_1$ <br> $h = \frac{m_2 + m_3}{m_1} > h_m$  Amplitude of<br>fundamental<br>- A reco **armonic Content of Steady State Wavefor**<br>vious non-sinusoidal waveform in steady state is<br>ficient but unnecessary condition of forced oscillation<br>rmonic content<br>Harmonic content higher than a given threshold is an indica
- Harmonic content
	- Harmonic content higher than a given threshold is an indicator of nonsinusoidal waveform, and forced oscillation

$$
= \frac{m_2 + m_3}{m_1} > h_m
$$
Ampli  
funda

 $m$   $\ldots$   $m_1$   $m_2$   $m_3$   $m_4$ Amplitude of

Amplitude of ith harmonic

– A recommended value of the threshold is 0.11, which is the harmonic content index for a triangle wave

#### **Simple but Practical**





#### **Examples**



h=0.34>0.11 Forced oscillation



WECC FO, 2015 From Dan





# **2. Features of Start Up Waveform**

• Different intrinsic system dampings result in different features of start up waveform



Start up: the stage when the amplitude increases





## **Features of Start Up Waveform**

• The envelope of start up waveform  $Ae^{\sigma t} + B$ 

> Natural oscillation  $A > 0, B = 0, \sigma > 0$ Forced oscillation  $A < 0, B = -A, \sigma < 0$





Е

# **Features of Start Up Waveform 11 Up Waveform**<br>
External  $Ae^{\sigma t} + B$ <br>  $Y_i = X_i - X_{i-1} = A(e^{\sigma T} - 1)e^{(i-1)\sigma T}$ <br>  $Y_i = \sigma T$ <br>  $S = \sigma T$ <br>  $R = \sigma T$ <br>  $S = \sigma T$ Up Waveform<br>
orm  $Ae^{\sigma t} + B$ <br>  $= X_i - X_{i-1} = A(e^{\sigma T} - 1)e^{(i-1)\sigma T}$ <br>  $= \sigma T$ <br>
ed oscillation **latures of Start Up Waveform**<br>
a of start up waveform  $Ae^{\pi} + B$ <br>
de is  $\sigma$ <br>
value  $X_i$ <br>
of peak-peak value  $Y_i = X_i - X_{i-1} = A(e^{\sigma T} - 1)e^{(i-1)\sigma T}$ <br>  $Z_i = \ln Y_i = (\ln(A(e^{\sigma T} - 1)) - \sigma T) + \sigma T \cdot i$ <br>
al oscillation, S<-e: forced oscillat **Solution And Algebra 10 May are the Common Section And Algebra 10 May be above that the problem of**  $Ae^{\alpha t} + B$ **<br>
is**  $\sigma$ **<br>
we**  $X_i$ **<br>
peak-peak value**  $Y_i = X_i - X_{i-1} = A(e^{\sigma T} - 1)e^{(i-1)\sigma T}$ **<br>**  $= \ln Y_i = (\ln(A(e^{\sigma T} - 1)) - \sigma T) + \sigma T \cdot i$ **<br>
fo**

- The envelope of start up waveform  $Ae^{\sigma t} + B$
- the difference is σ
- **Steps** 
	- Peak-peak value *Xi*
- increment of peak-peak value  $Y_i = X_i X_{i-1} = A(e^{\sigma T} 1)e^{(i-1)\sigma T}$  $(i-1)\sigma T$ 
	- logarithm  $Z_i = \ln Y_i = (\ln(A(e^{\sigma T} 1)) \sigma T) +$
	- Linear fitting to get the slope  $S = \sigma T$
	- S>ε: natural oscillation, S<-ε: forced oscillation





#### **Examples**





Measured S=0.14>0, Natural Simulated S=-0.06<0, Forced





# **3. Spectral Methods** (From Ruichao and Dan)

- The actual system response contains three components
- The intrinsic damping is contained in the ambient component







#### **Steady State Response**

• Natural (undamped)

$$
\hat{y}_r(t) = 2 \left| \underline{c}_r \underline{u}_n \underline{v}_n \underline{x}(0) \right| \cos \left( \omega_n t + \angle \left( \underline{c}_r \underline{u}_n \underline{v}_n \underline{x}(0) \right) \right) \Rightarrow \text{Transient}
$$
\n
$$
+ \sum_{l=1}^M \left[ q_l(t) \circledast \left[ 2 \left| \underline{c}_r \underline{u}_n \underline{v}_n \underline{b}_2 \right| \cos \left( \omega_n t + \angle \underline{c}_r \underline{u}_n \underline{v}_n \underline{b}_2 \right) \right] \right] \Rightarrow \text{Noise 1: Random Sinusoidal Noise}
$$
\n
$$
+ \sum_{l=1}^M \left[ q_l(t) \circledast \left[ \sum_{\substack{l=1 \ l \neq n, n^*}}^N \underline{c}_r \underline{u}_l \underline{v}_l \underline{b}_2 \right] e^{\lambda_l t} \right] \Rightarrow \text{Noise 2: Colored Noise}
$$

• **Forced**  
\n
$$
\hat{y}_r(t) = \sum_{m=1}^{\infty} \left( \left( \sum_{i=1}^N \underline{c}_r \underline{u}_i \underline{v}_i \underline{b}_1 e^{\lambda_i t} \right) \odot |A_m| \cos(m\omega_0 t + \angle A_m) \right) \Rightarrow \text{Forced}
$$
\n
$$
+ \sum_{l=1}^M \left( q_l(t) \odot \left( \sum_{i=1}^N \underline{c}_r \underline{u}_i \underline{v}_i \underline{b}_2 e^{\lambda_i t} \right) \right) \Rightarrow \text{Colored Noise}
$$





# **Signal-Noise Separation**

- First separate the whole response into signal and noise
- Power spectral density (PSD) of natural oscillation

 $S_{\hat{v}_r \hat{v}_r}(\omega_n) = S_{\hat{v}_{sr} \hat{v}_{sr}}(\omega_n) + S_{\hat{v}_{Nr} \hat{v}_{Nr}}(\omega_n)$ 

 $S_{\hat{x}_{sr}\hat{x}_{sr}}(\omega_n) = |2\pi \underline{c}_r \underline{u}_n \underline{v}_n \underline{x}(0)|^2 \delta(0)^2$ 

PSD of signal/transient component

$$
S_{\hat{y}_{Nr}\hat{y}_{Nr}}(\omega_n) \cong \left|2\pi\underline{c_r}\underline{u}_n\right|^2 \left[\sum_{l=1}^M \left| \underline{v_n}\underline{b}_{2l} \right|^2 S_{q_lq_l}(\omega) \right] \delta(0)^2
$$

PSD of noise, dominated by sinusoidal noise

• For two different measurements

$$
\alpha_{S} = \frac{S_{\hat{y}_{S1}\hat{y}_{S1}}(\omega_{n})}{S_{\hat{y}_{S2}\hat{y}_{S2}}(\omega_{n})} = \frac{|c_{1}u_{n}|^{2}}{|c_{2}u_{n}|^{2}} \quad \text{Equal} \quad \alpha_{N} = \frac{S_{\hat{y}_{N1}\hat{y}_{N1}}(\omega_{n})}{S_{\hat{y}_{N2}\hat{y}_{N2}}(\omega_{n})} = \frac{|c_{1}u_{n}|^{2}}{|c_{2}u_{n}|^{2}}
$$





## **Signal-Noise Separation**

2

 $\delta(0)^2$ 

• Power spectral density (PSD) of forced oscillation

 $\sum^N$  2π $A_1$ <u>c<sub>r</sub> u<sub>i</sub> v<sub>i</sub> b<sub>1</sub></u>

 $j\omega_0 - \lambda_i$ 

$$
S_{\hat{y}_r\hat{y}_r}(\omega_0) = S_{\hat{y}_{Sr}\hat{y}_{Sr}}(\omega_0) + S_{\hat{y}_{Nr}\hat{y}_{Nr}}(\omega_0)
$$

 $i=1$ 

PSD of signal/forced component

$$
S_{\hat{y}_{rn}\hat{y}_{rn}}(\omega_0) = \sum_{l=1}^{M} \left| \sum_{i=1}^{N} \frac{c_r u_i v_i b_{2l}}{j \omega_0 - \lambda_i} \right|^2 S_{q_l q_l}(\omega_0)
$$
 PSD of noise

• For two different measurements

 $S_{\hat{\mathcal{Y}}_{rf}\hat{\mathcal{Y}}_{rf}}(\omega_0) = \Big|\sum_{\mathcal{Y}}$ 

$$
\alpha_S = \frac{S_{\hat{y}_{S1}\hat{y}_{S1}}(\omega_n)}{S_{\hat{y}_{S2}\hat{y}_{S2}}(\omega_n)}
$$
 NOT Equal  $\alpha_N = \frac{S_{\hat{y}_{N1}\hat{y}_{N1}}(\omega_n)}{S_{\hat{y}_{N2}\hat{y}_{N2}}(\omega_n)}$ 





#### **Examples**

Natural  $\alpha_s$  = 3.80 and  $\alpha_N$  = 3.85



Power & Energy Society®

*Forced*  $\alpha_s$  = 7.28 and  $\alpha_N$  = 0.91





# **Cross-Spectrum Difference Function**

- To detect the existence of the "random sinusoid" component
- cross-spectrum difference function

 $S_r[\Omega] \triangleq \tilde{Y}_{rw_1}^* \tilde{Y}_{rw_2} - \tilde{Y}_{rw_2}^* \tilde{Y}_{rw_3}$ 

 $\tilde{Y}_{rw_i}$  the scaled DFT of the signal over window i

• Cross-spectrum index

$$
C_{rg}[\Omega] \triangleq \frac{\left| \mathbf{E}\left\{ S_{r}^{*}[\Omega]S_{g}[\Omega]\right\} \right|^{2}}{\mathbf{E}\left\{ S_{r}^{*}[\Omega]S_{r}[\Omega]\right\} \mathbf{E}\left\{ S_{g}^{*}[\Omega]S_{g}[\Omega]\right\}}
$$

r, g: different measurement channels

**Criterion** 





### **Examples**

Natural Forced



Power & Energy Society®





29

Е

### **Examples**

An actual forced oscillation incident in the western North American power system







# **Conclusions**

- $\triangleright$  Some methods for online distinguishing natural and forced oscillations are proposed
- $\triangleright$  The problem is NOT well solved
- $\triangleright$  Due to the complexity of actual oscillation curves, many methods, though have solid theoretical foundations and perform well with simulation results, do not perform well with actual records
- More practical approaches are still needed





# **Thanks!**

#### **Lei CHEN**

PhD, Associate Professor Department of Electrical Engineering Tsinghua University chenlei08@tsinghua.edu.cn



