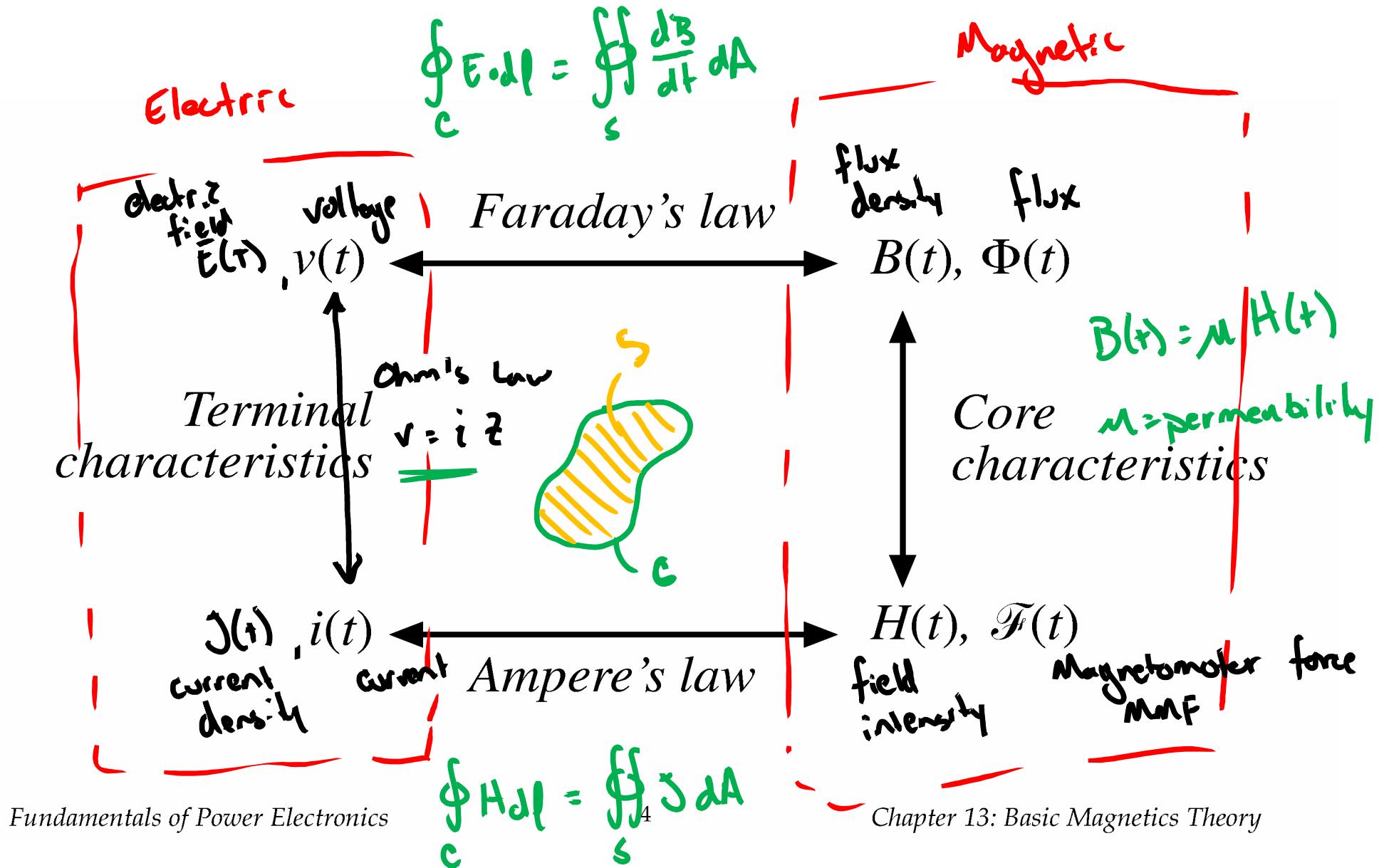
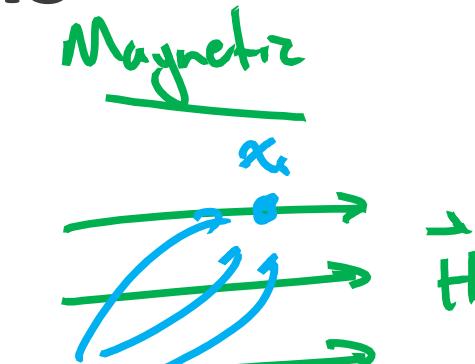
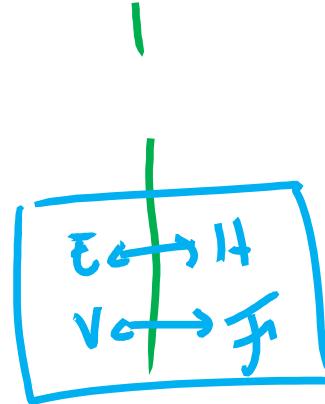
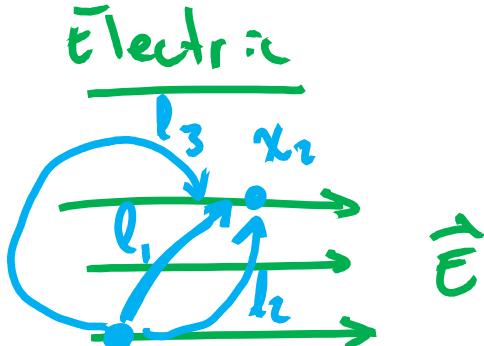


Basic Magnetics Relationships

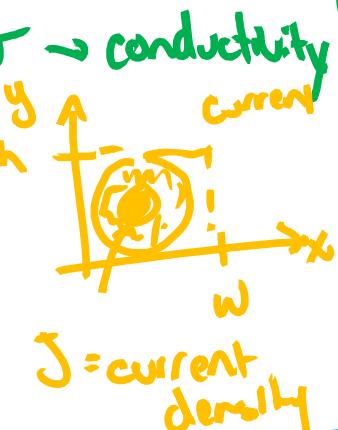
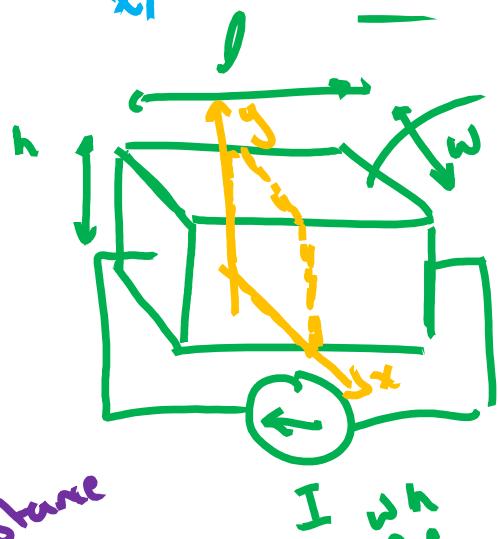


Electric/Magnetic Duals



$$\int_{x_1}^{x_2} \vec{E} dl = V_{x_2, x_1}$$

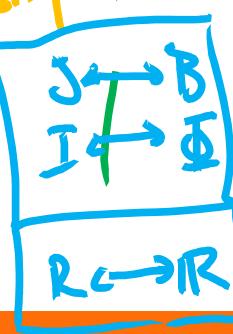
$$\int_{x_1}^{x_2} H dl = \mathcal{F}_{x_2, x_1}$$



resistance

$$R = \frac{1}{\sigma} \frac{l}{h \omega}$$

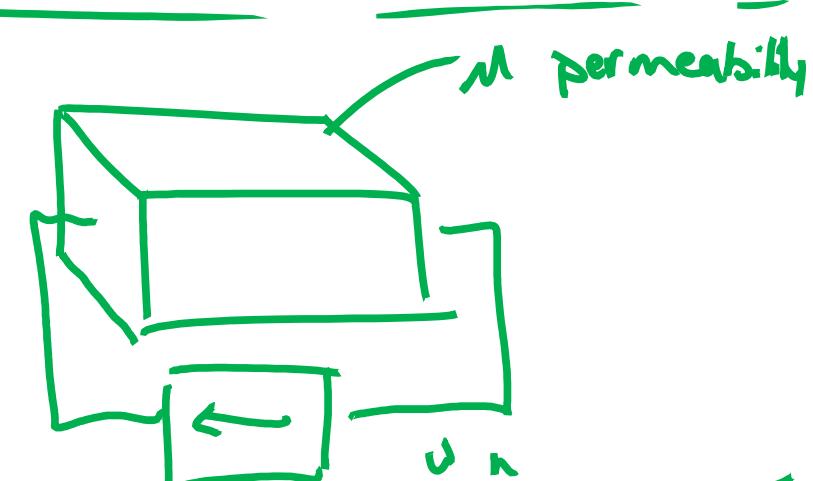
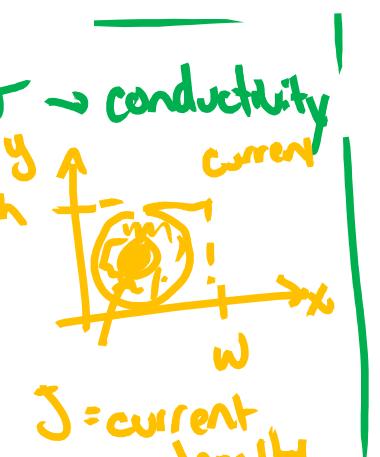
$$\iiint_{0}^h J dy dx = I$$



resistance

$$R = \frac{1}{\mu} \frac{l}{h \omega}$$

$$\iiint_{0}^h B dy dx = \Phi$$



Faraday's Law

$$V_{\text{turn}} = \oint_C E \cdot d\ell = \iint_S \frac{d\Phi}{dt} dA = \frac{d\Phi(t)}{dt} = \frac{dB(t)}{dt} A_c$$

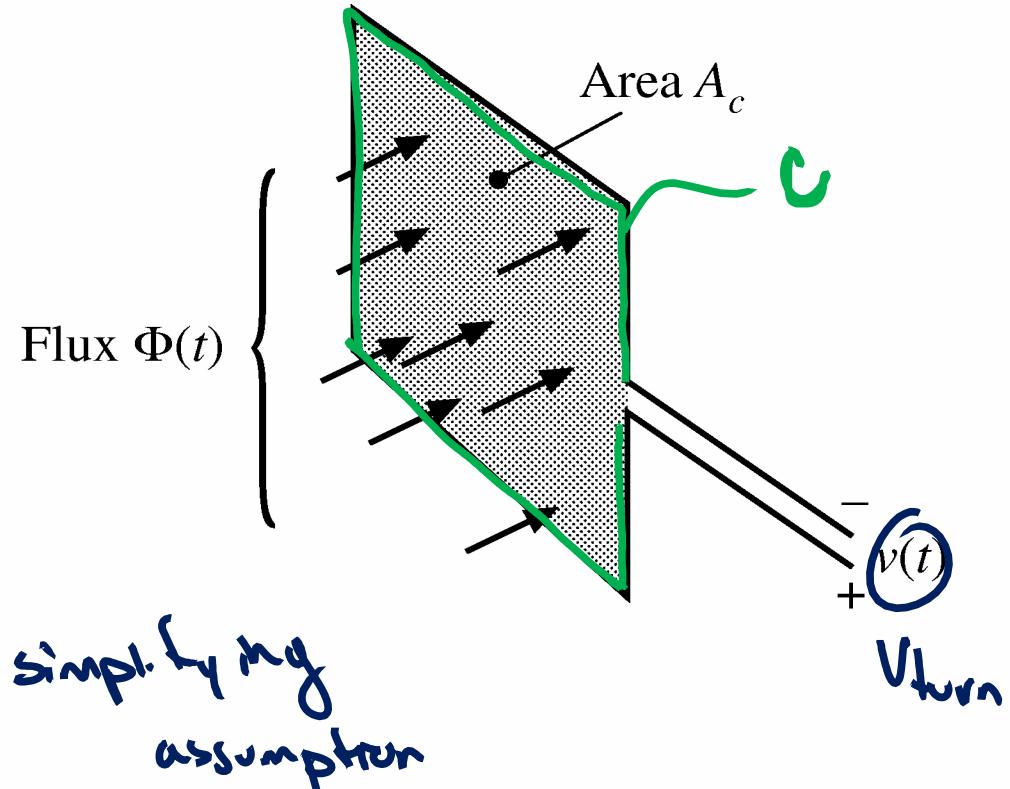
Voltage $v(t)$ is induced in a loop of wire by change in the total flux $\Phi(t)$ passing through the interior of the loop, according to

$$v(t) = \frac{d\Phi(t)}{dt}$$

For uniform flux distribution,
 $\Phi(t) = B(t)A_c$ and hence

$$v(t) = A_c \frac{dB(t)}{dt}$$

V_{turn}



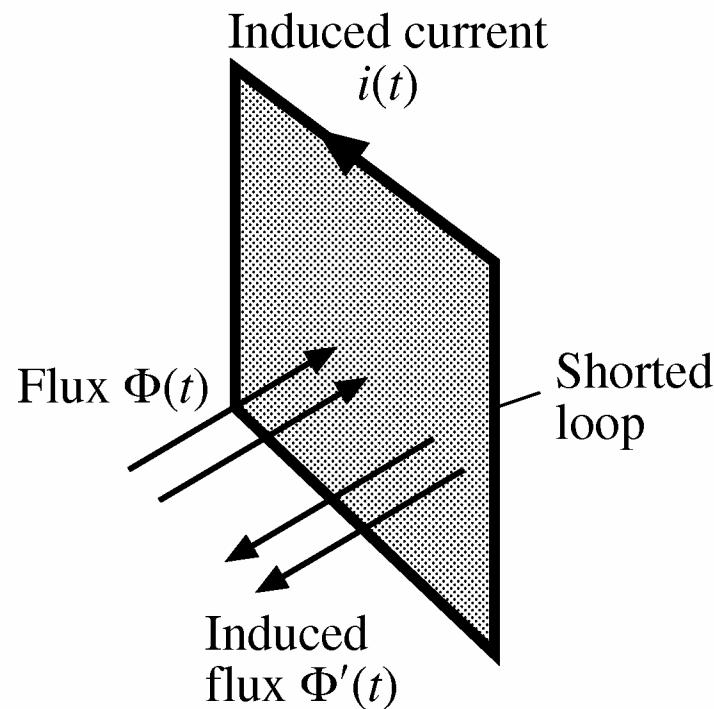
Lenz's Law

Right hand rule

The voltage $v(t)$ induced by the changing flux $\Phi(t)$ is of the polarity that tends to drive a current through the loop to counteract the flux change.

Example: a shorted loop of wire

- Changing flux $\Phi(t)$ induces a voltage $v(t)$ around the loop
- This voltage, divided by the impedance of the loop conductor, leads to current $i(t)$
- This current induces a flux $\Phi'(t)$, which tends to oppose changes in $\Phi(t)$



Ampere's Law

$$\mathcal{F} = H\ell_m = \oint_c \mathbf{H} \cdot d\ell = \iint_S \mathbf{B} \cdot d\mathbf{A} = I_{\text{enclosed}} = i(t)$$

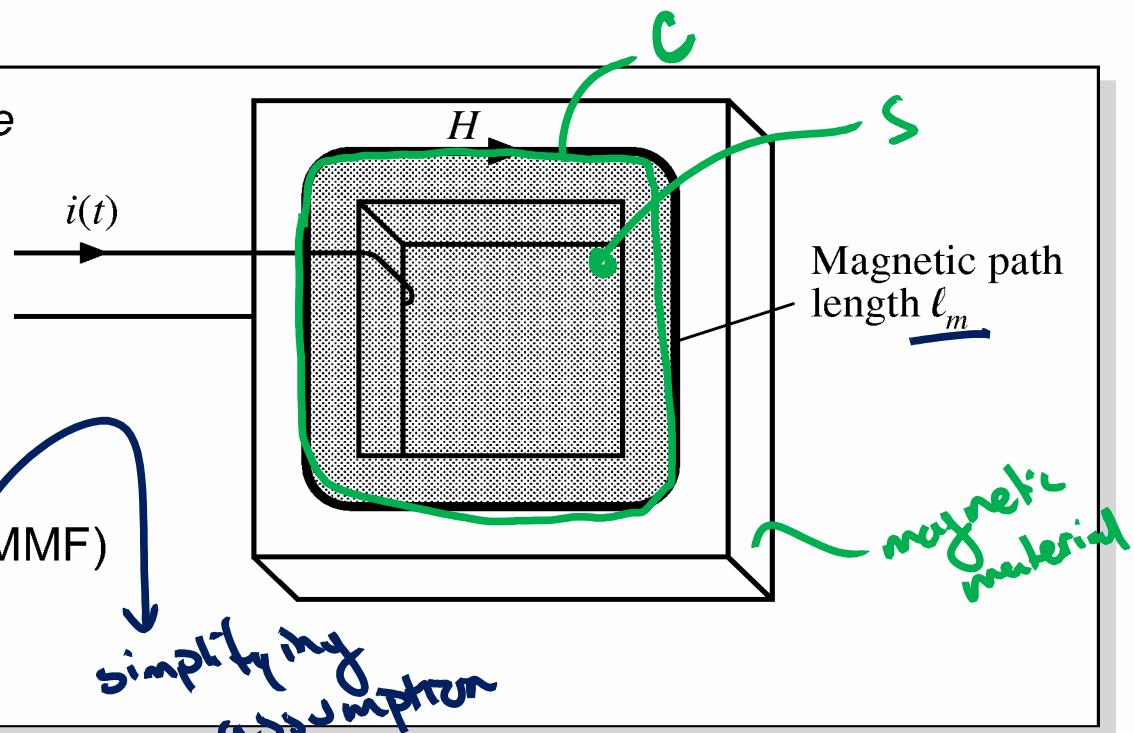
The net MMF around a closed path is equal to the total current passing through the interior of the path:

$$\oint_{\text{closed path}} \mathbf{H} \cdot d\ell = \text{total current passing through interior of path}$$

Example: magnetic core. Wire carrying current $i(t)$ passes through core window.

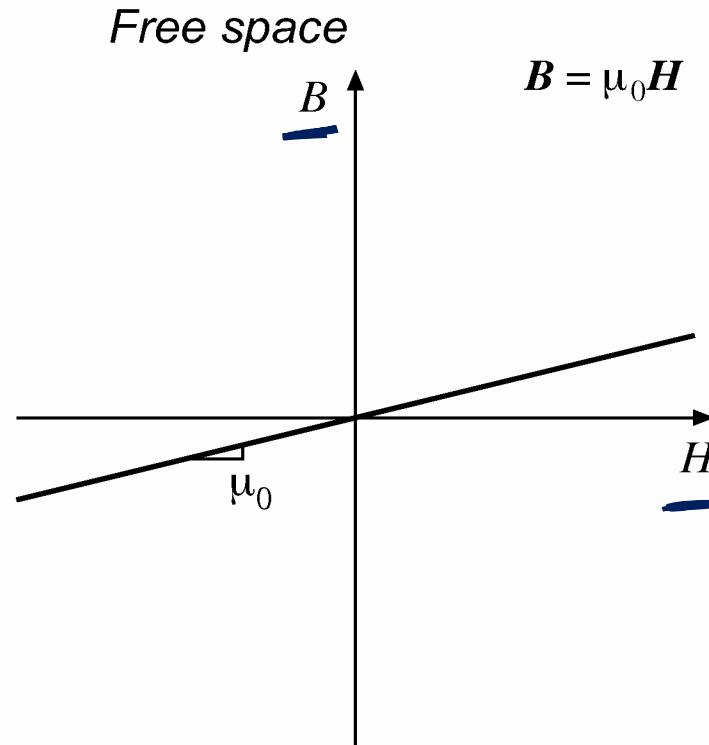
- Illustrated path follows magnetic flux lines around interior of core
- For uniform magnetic field strength $H(t)$, the integral (MMF) is $H(t)\ell_m$. So

$$\mathcal{F}(t) = H(t)\ell_m = i(t)$$



Core Material Characteristics

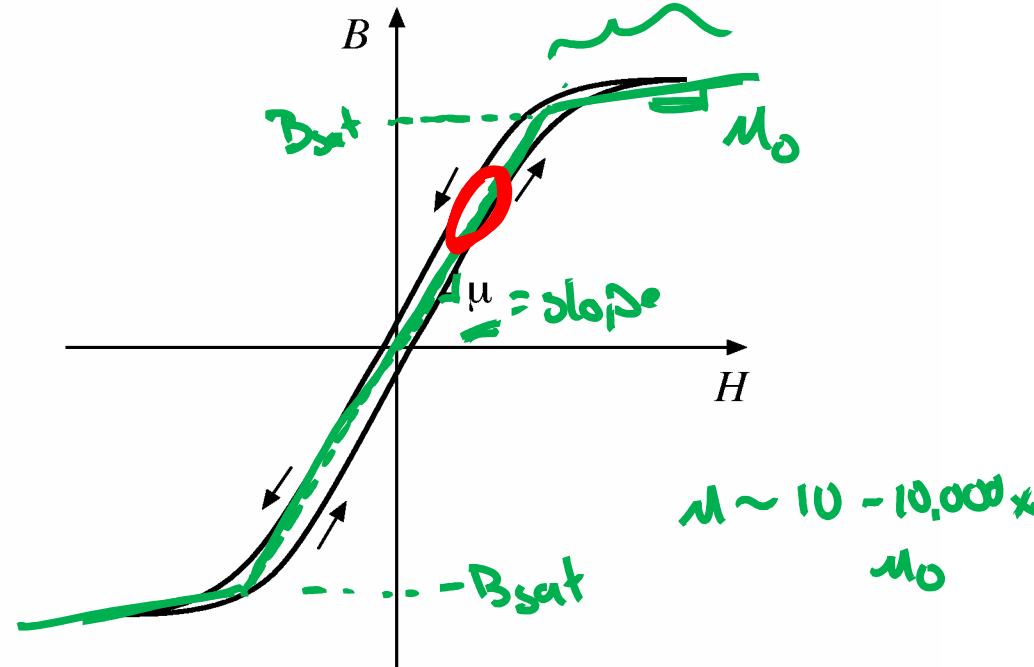
$$\underline{B = \mu H}$$



μ_0 = permeability of free space
= $4\pi \cdot 10^{-7}$ Henries per meter

Fundamentals of Power Electronics

A magnetic core material saturation



Highly nonlinear, with hysteresis and saturation

12

Chapter 13: Basic Magnetics Theory

Units

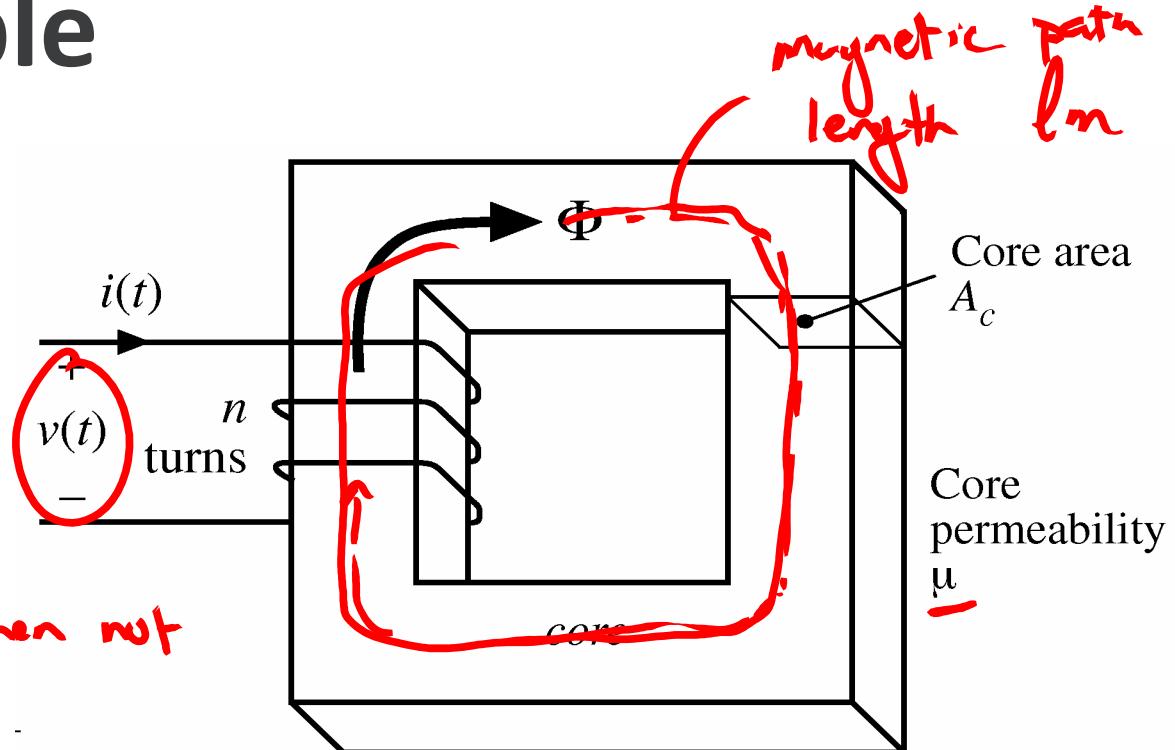
Table 12.1. Units for magnetic quantities

quantity	MKS	unrationalized cgs	conversions
core material equation	$B = \mu_0 \mu_r H$	$B = \mu_r H$	
B	Tesla	Gauss	$1\text{T} = 10^4\text{G}$
H	Ampere / meter	Oersted	$1\text{A/m} = 4\pi \cdot 10^{-3} \text{ Oe}$
Φ	Weber	Maxwell	$1\text{Wb} = 10^8 \text{ Mx}$ $1\text{T} = 1\text{Wb} / \text{m}^2$

Inductor Example

Simplifying assumptions:

- ① $M \gg \mu_0$ & flux stays entirely within the core
- ② H & B fields are uniform throughout the core
- ③ $B = \mu H$ in the core (when not saturated)



Faraday:

$$V_{\text{turn}} = A_c \frac{dB(t)}{dt}$$

$$V(t) = n A_c \frac{dB(t)}{dt}$$

Ampere:

$$H l_m = I_{\text{enclosed}}$$

$$H l_m = n i(t)$$

Core Material:

$$B = \mu H \quad \text{if } |B| < B_{\text{sat}}$$

$$V(t) = n A_c \frac{d}{dt} \left(\mu \frac{n i(t)}{l_m} \right) = \frac{n^2 A_c \mu}{l_m} \frac{di(t)}{dt}$$